Quantum Computation Estimating eigenvalues: An application of QFT

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The problem: Eigenvalue estimation

Several algorithms previously discussed (Simon, Deutsch-Joza, etc) resort to the following technique:

- Take a controlled version of an operator *U* and prepare the target qubit with an eigenvector,
- win order to push up (or kick back) the associated eigenvalue to the state of the control qubit as in

$$cU(a_0|0\rangle+a_1|1\rangle)\left(\frac{|0\rangle-|1\rangle}{\sqrt{2}}\right) = \left((-1)^{f(0)}a_0|0\rangle+(-1)^{f(1)}a_1|1\rangle\right)\left(\frac{|0\rangle-|1\rangle}{\sqrt{2}}\right)$$

When the eigenvector is difficult to build $_{\rm OOO}$

The problem: Eigenvalue estimation

The question

Can this technique be generalised to estimate the eigenvalues of an arbitrary, *n*-qubit unitary operator U?

The eigenvalue estimation problem

Let $(|\psi\rangle, e^{i2\pi\phi})$, with $0 \le \phi < 1$, be an eigenvector, eigenvalue pair for a unitary U. Determine ϕ .

Note that eigenvalues of unitary operators are always of this form.

Quantum eigenvalue estimation

Algorithm performance

When the eigenvector is difficult to build $_{\rm OOO}$

The strategy

Use a controlld version of U to prepare a state from which ϕ can be found through the inverse of the QFT (recall the phase estimation problem already discussed)

A Simple Example

Take a unitary U with an eigenvector $|\psi\rangle$ whose eigenvalue is $e^{i2\pi\phi} \phi$ is equal to one of the values $\{0 \cdot \frac{1}{2}, 1 \cdot \frac{1}{2}\}$. Find out ϕ .

This is obtained via the circuit



When the eigenvector is difficult to build $_{\rm OOO}$

The general case

In less trivial cases, a multi-controlled version of U is reguired:

$$\begin{bmatrix} \swarrow^{n} & \uparrow^{n} \\ \swarrow^{m} & U & \swarrow^{m} \end{bmatrix} = |x\rangle |y\rangle \mapsto |x\rangle U^{x} |y\rangle$$

Intuitively it applies U to $|y\rangle$ a number of times equal to x

Examples

 $\ket{10}\ket{y}\mapsto \ket{10}\left(UU\ket{y}
ight)$ and $\ket{00}\ket{y}\mapsto \ket{00}\ket{y}$

Note that $|\psi\rangle$ is also an eigenvector of U^x , with eigenvalue $e^{i2\pi x\phi}$, for any integer x.

Multi-controlled operations

Recall that a binary number $x_1 \dots x_n$ corresponds to the natural number $2^{n-1}x_1 + \dots + 2^0x_n$

We use this to build the previous multi-controlled operation in terms of simpler operations



Note that the multi-controlled operation uses n 'simply'-controlled rotations U^{2^i}

When the eigenvector is difficult to build $_{\rm OOO}$

Another Example

Take a unitary U with an eigenvector $|\psi\rangle$ whose eigenvalue is $e^{i2\pi\phi}$ ϕ is equal to one of the following values $\left\{0 \cdot \frac{1}{4}, 1 \cdot \frac{1}{4}, 2 \cdot \frac{1}{4}, 3 \cdot \frac{1}{4}\right\}$

The following circuit discovers ϕ



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Another Example

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The following circuit discovers ϕ



 $\begin{array}{c} \mbox{Quantum eigenvalue estimation} \\ \mbox{000000000} \end{array} \\ \end{array}$

Algorithm performance

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Another Example



$$\begin{split} |0\rangle |0\rangle \\ \stackrel{H^{\otimes 2}}{\mapsto} \frac{1}{\sqrt{2^2}} (|00\rangle + |01\rangle + |10\rangle + |11\rangle) \\ \stackrel{\text{ctrl. }U}{\mapsto} \frac{1}{\sqrt{2^2}} (|00\rangle + e^{i2\pi\phi} |01\rangle + e^{i2\pi\phi\cdot 2} |10\rangle + e^{i2\pi\phi\cdot 3} |11\rangle) \\ = \frac{1}{\sqrt{2^2}} (|00\rangle + e^{i2\pi\times\cdot\frac{1}{4}} |01\rangle + e^{i2\pi\times\cdot\frac{1}{4}\cdot 2} |10\rangle + e^{i2\pi\times\cdot\frac{1}{4}\cdot 3} |11\rangle) \\ = \frac{1}{\sqrt{2^2}} (|00\rangle + \omega_2^{x} |01\rangle + \omega_2^{x\cdot 2} |10\rangle + \omega_2^{x\cdot 3} |11\rangle) \\ \stackrel{QFT_2^{-1}}{\mapsto} |x\rangle \end{split}$$

When the eigenvector is difficult to build $_{\rm OOO}$

Yet Another Example

Take a unitary U with eigenvector $|\psi\rangle$ whose eigenvalue is $e^{i2\pi\phi}$ st $\phi\in \left\{0\cdot\frac{1}{2^n},\ldots,2^n-1\cdot\frac{1}{2^n}\right\}$

The following circuit returns x such that $\phi = x \cdot \frac{1}{2^n}$



Exercise

Prove that indeed the circuit returns x such that $\phi = x \cdot \frac{1}{2^n}$

Quantum eigenvalue estimation

Algorithm performance

When the eigenvector is difficult to build $_{\rm OOO}$

Yet Another Example

Exercise

Show that $QFT_n |0\rangle = H^{\otimes n} |0\rangle$.

Note that this allows to rewrite the previous circuit in the one below



encodes x in local phases (in the form of rotations)

When the eigenvector is difficult to build $_{\rm OOO}$

Complexity of quantum eigenvalue estimation



encodes k in local phases (in the form of rotations)

How many gates does the circuit above use?

n 'Hadamards' + *n* 'simply'-controlled gates + n^2 gates for QFT_n^{-1}

When the eigenvector is difficult to build $_{\rm OOO}$

... but precision is Limited

We assumed $0 \le \phi < 1$ takes a value from $\left\{0 \cdot \frac{1}{2^n}, \dots, 2^n - 1 \cdot \frac{1}{2^n}\right\}$... an assumption that arose from having only *n* qubits to estimate ...

But what to do if ϕ takes none of these values? Return the *n*-bit number *k* with $k \cdot \frac{1}{2^n}$ the value above closest to ϕ

Is the circuit above up to this task?

When the eigenvector is difficult to build $_{\rm OOO}$

Setting the stage

Let
$$\omega_n = e^{i2\pi \cdot \frac{1}{2^n}}$$
 (division of the unit circle in 2^n slices)

To answer the previous question, we will use the following explicit defn. of QFT^{-1}

$$QFT_n^{-1} \left| x \right\rangle = \frac{1}{\sqrt{2^n}} \sum_{k=0}^{2^n - 1} \omega_n^{-x \cdot k} \left| k \right\rangle$$

Exercise

Prove that QFT_n^{-1} is indeed the inverse of QFT_n

When the eigenvector is difficult to build $_{\rm OOO}$

Setting the stage

Let
$$k \cdot \frac{1}{2^n}$$
 be the value in $\left\{0 \cdot \frac{1}{2^n}, \dots, 2^n - 1 \cdot \frac{1}{2^n}\right\}$ closest to ϕ , i.e.
 $\exists_{\epsilon} \cdot 0 \leq |\epsilon| \leq \frac{1}{2^n}$ and $k \cdot \frac{1}{2^n} + \epsilon = \phi$

Note that the difference ϵ decreases when the number of qubits increases.

Recall the circuit



encodes k in local phases (in the form of rotations)

When the eigenvector is difficult to build $_{\rm OOO}$

Setting the stage

Let
$$k \cdot \frac{1}{2^n}$$
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Computing the output again

$$\begin{split} |0\rangle \\ \stackrel{H^{\otimes n}}{\mapsto} & \frac{1}{\sqrt{2^{n}}} (|0\rangle + |1\rangle + \dots + |2^{n} - 1\rangle) \\ \stackrel{\text{ctrl. }U}{\mapsto} & \frac{1}{\sqrt{2^{n}}} \Big(|0\rangle + e^{i2\pi\phi \cdot 1} |1\rangle + \dots + e^{i2\pi\phi \cdot 2^{n-1}} |2^{n} - 1\rangle \Big) \\ &= \frac{1}{\sqrt{2^{n}}} \Big(|0\rangle + e^{i2\pi(k \cdot \frac{1}{2^{n}} + \epsilon) \cdot 1} |1\rangle + \dots + e^{i2\pi(k \cdot \frac{1}{2^{n}} + \epsilon) \cdot 2^{n-1}} |2^{n} - 1\rangle \Big) \\ &= \frac{1}{\sqrt{2^{n}}} \sum_{j=0}^{2^{n}-1} e^{i2\pi(k \cdot \frac{1}{2^{n}} + \epsilon) \cdot j} |j\rangle \\ &= \frac{1}{\sqrt{2^{n}}} \sum_{j=0}^{2^{n}-1} e^{i2\pi k \cdot \frac{1}{2^{n}} \cdot j} e^{i2\pi \epsilon \cdot j} |j\rangle \\ \stackrel{QFT^{-1}}{\mapsto} & \frac{1}{\sqrt{2^{n}}} \sum_{j=0}^{2^{n}-1} e^{i2\pi k \cdot \frac{1}{2^{n}} \cdot j} e^{i2\pi \epsilon \cdot j} \Big(\frac{1}{\sqrt{2^{n}}} \sum_{l=0}^{2^{n}-1} e^{-i2\pi j \cdot \frac{1}{2^{n}} \cdot l} |l\rangle \Big) \\ &= \frac{1}{2^{n}} \sum_{j=0}^{2^{n}-1} e^{i2\pi k \cdot \frac{1}{2^{n}} \cdot j} e^{i2\pi \epsilon \cdot j} \Big(\sum_{l=0}^{2^{n}-1} e^{-i2\pi j \cdot \frac{1}{2^{n}} \cdot l} |l\rangle \Big) \\ &= \frac{1}{2^{n}} \sum_{j=0}^{2^{n}-1} \sum_{l=0}^{2^{n}-1} e^{i2\pi \epsilon \cdot j} e^{i2\pi \epsilon \cdot j} (k-l) |l\rangle \end{split}$$

When the eigenvector is difficult to build $_{\rm OOO}$

Looking into the final state

The amplitude of $|k\rangle$ is $\frac{1}{2^n} \sum_{j=0}^{2^n-1} e^{i2\pi\epsilon \cdot j}$ which is a finite geometric series. Therefore,

$$\frac{1}{2^n}\sum_{j=0}^{2^n-1}e^{i2\pi\epsilon j} = \begin{cases} 1 & \text{if } \epsilon = 0\\ \frac{1}{2^n}\frac{1-e^{i2\pi\epsilon 2^n}}{1-e^{i2\pi\epsilon}} & \text{if } \epsilon \neq 0 \end{cases}$$

Let us proceed under the assumption $\epsilon \neq 0$.

When the eigenvector is difficult to build $_{\rm OOO}$

A geometric detour

 $|1 - e^{i\theta}|$ for some angle θ is the Euclidean distance between 1 and $e^{i\theta}$ (length of the straight line segment between both points) Consider also arc length θ between 1 and $e^{i\theta}$ (distance between the two points by running along the unit circle)

Theorem

Let d^E and d^a be respectively the Euclidean distance and arc length between 1 and $e^{i\theta}$. Then,

a.
$$d^{E} \leq d^{a}$$

b. if $0 \leq \theta \leq \pi$ we have $\frac{d^{a}}{d^{E}} \leq \frac{\pi}{2}$

When the eigenvector is difficult to build 000

Finally!

Recall
$$\left|\frac{1}{2^{n}}\frac{1-e^{i2\pi\epsilon^{2^{n}}}}{1-e^{i2\pi\epsilon}}\right|^{2}$$
 is the probability of measuring $|k\rangle$
 $\left|\frac{1}{2^{n}}\frac{1-e^{i2\pi\epsilon^{2^{n}}}}{1-e^{i2\pi\epsilon^{2^{n}}}}\right|^{2} = \left(\frac{1}{2^{n}}\right)^{2}\frac{\left|1-e^{i2\pi\epsilon^{2^{n}}}\right|^{2}}{\left|1-e^{i2\pi\epsilon^{2^{n}}}\right|^{2}}$ {Thm. [a.]}
 $\geq \left(\frac{1}{2^{n}}\right)^{2}\frac{\left(\frac{1}{2^{n}} \cdot 2\pi\epsilon^{2^{n}}\right)^{2}}{(2\pi\epsilon)^{2}}$ {Thm. [b.]}
 $= \left(\frac{1}{2^{n}}\right)^{2}\frac{\left(4\epsilon^{2^{n}}\right)^{2}}{(2\pi\epsilon)^{2}}$ {Thm. [b.]}

Working with a superposition of eigenvectors

The algorithm requires an eigenvector as input, but sometimes is highly difficult to build such a vector.

Often it is easier to feed instead a superposition of eigenvectors.

Indeed, by the spectral theorem one knows that the eigenvectors $\{|\psi_1\rangle, \ldots, |\psi_N\rangle\}$ of U (with associated eigenvalues $e^{i2\pi\phi_1}, \ldots, e^{i2\pi\phi_N}$) form a basis for the $N(=2^n)$ -dimensional vector space on which U acts.

Thus, one may define

$$|\psi\rangle = \frac{1}{\sqrt{N}}(|\psi_1\rangle + \cdots + |\psi_N\rangle)$$

to feed the circuit

When the eigenvector is difficult to build $\circ \bullet \circ$

Working with a superposition of eigenvectors



encodes k in local phases (in the form of rotations)

Exercise

Show that if $\forall_{i \leq N} \cdot \phi_i \in \left\{ 0 \cdot \frac{1}{2^n}, \dots, 2^n - 1 \cdot \frac{1}{2^n} \right\}$ then the circuit's output is

$$\frac{1}{\sqrt{N}} \left(\left| k_1 \right\rangle \left| \psi_1 \right\rangle + \dots + \left| k_N \right\rangle \left| \psi_N \right\rangle \right) \qquad \left(\phi_i = k_i \cdot \frac{1}{2^n} \right)$$

Working with a superposition of eigenvectors

Similarly, it can be shown that in general the circuit's output is

$$\frac{1}{\sqrt{N}}\left(\left|\tilde{\phi_{1}}\right\rangle\left|\psi_{1}\right\rangle+\cdots+\left|\tilde{\phi_{N}}\right\rangle\left|\psi_{N}\right\rangle\right)$$

where each $\tilde{\phi}_i$ is the best *n*-bit approximation of ϕ_i with probability $\geq \frac{4}{\pi^2}$